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Analytical Solution of Thick Rectangular Plate with Clamped and Free Support Boundary Condition using Polynomial Shear Deformation Theory

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ARTICLE INFO ABSTRACT

In this paper, a new polynomial shear deformation theory for the static flexural analysis of anisotropic rectangular thick plate was developed. The plate which carries a uniformly distributed load is clamped on the three edges, and free of support on the other edge (CCFC), is analyzed to determine the in-plane displacement, vertical displacement, bending moment, and shear force, bending and transverse shear stress. The General *variation approach was used to obtain the general governing equation and its associated boundary conditions, thereafter the coefficient of deflection and shear deformation along the direction of x and y coordinate was determined by minimizing the energy equation obtained using the new established theory. The study revealed that: (i) as the displacement and stress decrease, the plate's span-thickness ratio increases (ii) as the length to breadth ratio of the plate increases, the value of displacement and stresses increase. To validate the theory, the numerical results are obtained and compared with an available solution in the literature. The result showed good agreement with those in the literature.*

1. Introduction

Plates, generally well-known structural element have been extensively used in different fields like engineering, and other industries due to their high strength, lightweight, load resistance characteristic and their cost benefits, as recorded in [1]. Plates are widely used in structural engineering, aerospace engineering, mechanical engineering, etc., for the construction of buildings, retaining walls, ships, bridges, railways, etc., [2 - 3].

Plates have been classified based on thickness as: thick, moderately thick, and thin plates [4], based on material composition as: isotropic and anisotropic plates, based on shape as: circular, rectangular, triangular, elliptical, and skew plates, [5]. The edges of plates can be of different kinds of supports such as: simply supported support, point and fixed supports, etc. as seen in [6].

The analysis or solutions of plates help to determine the stability of plates, and it entails the determination of

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displacements, stresses, moments, etc., at a different portion of the plate, [7]. Also, the analysis of plates can be due to bending, buckling, and vibration caused forces. It is generally known that the bending of plates relates to the deflection of the plate at a right angle (90°) to the plate plane due to the influence of forces and moments [8].

It is the depth of plates that mainly affects their bending characteristics compared to other surface dimensions like: length and breadth [9].

There have been different theorems: the classical plate theory (CPT) and refined plate theory (RPT) developed and applied in the analysis of the different categories of plates by many researchers. The classical plate theory (CPT), which was developed by [10 - 11], is based on the assumption that the line which is normal to the mid-surface remains so, before and after deformation. The classical plate theory (CPT) is mainly used in the analysis of thin plates and has also been used by [12 - 14] in the analysis of solutions involving plates and shells, respectively. But it has been found to be inadequate by [15] for the analysis of thick plates as its neglects the effect of shear deformation $-$ a gap that led to the development of refined plate theory by researchers.

The Refined Plate Theory (RPT) is used for the analysis of thick plates [16], and it addresses the gap in classical plate theory (CPT).

The RPT comprises First-Order Shear Deformation Theory (FSDT – which has a correction factor as a limitation, and Higher-Order Shear Deformation Theories (HSDT), etc., taking into account the effect of shear deformation and has been applied by different researchers in analyzing thick plates. Authors in [17 - 18] revealed that the limitations in CPT and FSDTs led to the formulation of higher-order shear deformation theory (HSDT). There have been various studies on the bending analysis of thick rectangular plates with different boundary conditions using refined plated theory with different methods and techniques.

This present paper is aimed at bridging the gap in the literature by applying the polynomial shear deformation theory in obtaining the displacements, bending moment, and stresses of the thick rectangular plate. This theory, which is based on a traditional fourth-order shear deformation plate, is presented and applied in a bending analysis of thick rectangular plates using the Energyvariational approach. The investigation has conducted on the plate that has three edges clamped and the other opposite edge free of support (CCFC). Furthermore, derivation of general governing equations for the plate and numerical solutions for deformation, moment, and stress distributions of different points of the plate with a uniformly distributed transverse load is presented.

The physical interpretation of CCFC plate is that, the three clamped edges are supported by a beam and continuous over the span of the plate while the remaining one is free of support i.e. hang without support (Eg. Cantilever). This makes the study very significant because such BC exist depending on the type of beam/column support in the plate structure i.e. it contains different support case of the plate structure than other BCs. So, whenever such case (CCFC) as explained occurs in the structure, analyzing the plate as if is any other BC like SSSS, CCCC etc.(22, 27, 28 and 29) will not account for all the forces (stresses) acting on it. This is because, forces are generated due to the applied load on the structure thereby will introduce significant errors in the analysis and not predict an accurate or reliable result for the design if they are not considered. Thus, for a safe structural design of plate (slab) in a building or any type of structure with CCFC BCs, plate analysis with CCFC BCs like this is required.

The aim of this study is to determine the bending analysis of an isotropic rectangular plate for the effects of aspect ratio and displacement on the moment, shear force, stresses, and stress resultant of a thick rectangular plate. The study sought to achieve the aim through the following objectives:

- To generate the potential energy of a thick plate in line with the work of authors in [5].
- To formulate the general governing equations of the plate and obtain equations for the coefficients of deflection and shear deformation slope for x and y coordinates.
- To determine the expressions for the in-plane and out of plane displacement, bending moment, shear force and stresses of the thick rectangular plate.

2. Previous Works

The authors in [19] used Fourier trigonometric series to analyze the bending of rectangular plates on elastic foundations. They developed the deflection equation, which was in the form of trigonometric series using the Navier solution to the plate. The constraint of the trigonometric series is that it is tedious and rigorous for the formulation of shape functions for the analysis of rectangular plates. The authors in [20] stated that the constraint of the trigonometric series is that it is very strenuous for the formulation of shape functions relating to plates with unsymmetrical boundary conditions.

The authors in [21] used Mindlin's plate higher-order shear deformation theory to generate governing equations for the bending analysis of thick rectangular plates with clamped supports at all its edges (CCCC). Their developed theory involved the decoupling method and improvement on Navier's solutions with partial differential equations, which reduced the complexity involved in obtaining coefficients of plates. In as much as their results showed great accuracy when compared with past results of other studies, their study did not address plate with the CCFC edge condition and didn't make use of the direct variation energy method, which is less rigorous.

In [16], the authors used polynomial shape function for the analysis of rectangular isotropic thick plates. They used the direct variation energy method to obtain displacement coefficients without the need for correction factor for plates with CCCS and SSFS edge conditions, respectively. Their results (displacements and stresses), which they found to be adequate after comparing with results from previous works but did not include the CCFC plate's edge condition as it was not included in their study.

In [17], the authors used exponential function in shear deformation to determine the bending solution of the thick isotropic square and rectangular plates with simply supported supports. They applied the refined shear deformation theory and exponential functions in the form of thickness coordinate to obtain transverse shear deformation and rotary inertia without the need for a correction factor. They found their results: displacements, stresses, and frequencies to be satisfactory after comparing with other refined and exact theories. They only considered the boundary condition (SSSS) using the exponential function and not the polynomial displacement function in an energy method, which is less cumbersome.

The authors in [22] worked on the bending solutions of a thick rectangular plate with CCCC edge conditions. They used refined plate theory (RPT) - polynomial shear deformation theory (PSDT) with the Ritz energy method in their analysis without considering correction factor, and they obtained the coefficients for displacement by putting the total static energy equations to direct variation. They found out that their results (displacements and stresses) were accurate after comparing with results from other previous studies. The authors in [23] used the same boundary condition to study the bending behavior of the elastic rectangular thick plate using the Bergan-Wang approach. They found out that their approach can be used to determine the deflection of plates and analysis of thin, moderately thin, and thick plates, respectively. The authors in [22] used an assumed shape

function, which is not a close-form solution, while the authors in [23] didn't use the Energy method, which is more simplified. Both authors did not consider the plate with the CCFC edge condition in their analysis.

In [24], the authors did a study on bending solutions of the thick plate using a numerical method - based on the 3-D theory of elasticity. They used the method: spline collocation in two coordinate directions for their analysis. They concluded their approach to be useful in determining the distribution of displacements and stresses in thick rectangular plates, with hinged or clamped edge conditions. Their study did not cover the plate with the CCFC edge condition, and also they made use of a numerical method which, unlike the analytical approach, cannot determine the value of deflection at any given point in the plate.

The author in [25] used HSDT with stretching effect to determine bending characteristics of exponentially graded material plates. In his study, he made the assumption that the modulus of elasticity varies exponentially through the plate thickness direction. He found that his theory satisfied the shear stress-free condition at the top and bottom surface of the plates without considering the correction factor. He found out from his study that his theory was satisfactory after compared with previous studies and then concluded that the theory is suitable for predicting bending behavior of exponentially graded material plates. It can be seen from his study that he used shear strain shape function – a new higher-order shear deformation theory and not a polynomial displacement function in the energy method.

The authors in [26] have used PSDT for the analysis of rectangular plates. They obtained the expression for the critical lateral load that induces deflection and shear as a solution to the bending problem of rectangular plates with all four edges clamped (CCCC) and plate with free of support at the third edge and the other edges clamped (CCFC) using third-order shear deformation theory for thick rectangular plate respectively. They did not analyze the effect of displacement, moment, and stresses in predicting responses of the applied load on the plate structure.

Apart from the distinctiveness of the present study with respective individual previous works, there exists an aspect of the distinctiveness of the present study over the previous works put together. This lies in the type of shape functions, and plates support boundary conditions. Unlike the previous works that assumed the displacement function, the present work obtains the exact formulation of the total potential energy from a principle of static elastic theory to get a close form solution of the polynomial displacement function. Except for the author in [26], none of the researchers could determine the analytical solution of bending moment and stresses used in predicting the bending behavior of an elastic thick rectangular CCFC plate. These aspects cannot be overlooked.

3. Methodology

3.1. Assumptions

Considering the following assumptions, the total potential energy of a thick rectangular plate will be formulated. They include:

- The material of the plate is elastic and homogeneous in nature.
- \triangleright The vertical line that is initially normal to the middle surface of the plate before bending is neither straight nor normal to the middle surface after bending.
- \triangleright The effect of the strain and stress normal to the x-y plane on the gross response of the plate is small and can be neglected. Thus, the in-plane displacements, u and v, are differentiable in x, y, and z coordinates, while the deflection, w, is only differentiable in x and y coordinates.

3.2. Displacement, Strain and Stress Relationships

Figure 1 shown the two in-plane displacements: in-plane displacement in the direction of x coordinates (u) and in-plane displacement in the direction of y coordinates (v) are represented in Equations (1) and (2) (see the author in $[1]$):

Figure 1: Deformation of a section of a thick Plate

$$
u = \frac{z d \theta}{dx} + S(z) . \Omega_x \tag{1}
$$

$$
v = \frac{zd \cup}{dy} + S(z) . \cap y \tag{2}
$$

The shear deformation profile of plate section $S(z)$ is given (see the author in [5]) as:

$$
S(z) = \frac{5z}{3} \left(z^3 - \frac{2z}{t^3} \right)
$$
 (3)

where;

 \cap_x and \cap_y is the shear deformation slope in x and y axis respectively.

It is assume from this work that the out-of-plane displacement (∪) is only differentiable in x and y coordinates.ie. (∪=∪ (x, y) .

The normal stress along x-axes (ε_x) becomes:

$$
\varepsilon_x = \frac{du}{dx} \tag{4}
$$

The normal stress along y-axes (ε_v) gives:

$$
\varepsilon_y = \frac{du}{dy} \tag{5}
$$

The shear strain along $(x-y)$, $(x-z)$ and $(y-z)$ respectively are given in the Equation (6) , (7) and (8) respectively as:

$$
\gamma_{xy} = \frac{du}{dy} + \frac{dv}{dx} \tag{6}
$$

$$
\gamma_{xz} = \frac{du}{dz} + \frac{dU}{dx} \tag{7}
$$

$$
\gamma_{yz} = \frac{dv}{dz} + \frac{dU}{dx} \tag{8}
$$

dz dx
Similarly, the stress normal to x-axes is defined as:

$$
\sigma_x = \frac{E(\varepsilon_x + \mu \varepsilon_y)}{1 - \mu^2} \tag{9}
$$

the stress normal to y-axes is defined as:

$$
\sigma_y = \frac{E(\varepsilon_y + \mu \varepsilon_x)}{1 - \mu^2} \tag{10}
$$

Considering the relationship between the strain and stress in the plate, the shear strain in x-y plane is given as:

$$
\tau_{xy} = \frac{E}{2(1+\mu)} \cdot \frac{du}{dy} + \frac{dv}{dx}
$$
\n(11)

the shear strain in y-z plane is given as:

$$
\tau_{xz} = \frac{E}{2(1+\mu)} \cdot \frac{du}{dz} + \frac{dU}{dx}
$$
\nthe shear strain in y-z plane is given as:

\n
$$
\int_{E} \frac{du}{dx} \cdot \int_{U} \frac{dU}{dx} \cdot dU
$$
\n(12)

$$
\tau_{yz} = \frac{E}{2(1+\mu)} \cdot \frac{dv}{dz} + \frac{dU}{dx}
$$
\n(13)

Substituting Equation (1), (2), (4) and (5) into Equation (9), we have:

$$
\sigma_x = \frac{E}{1 - \mu^2} \left[\left(-\frac{z d^2}{dx^2} + \frac{S(z) d \Omega_x}{dx} \right) - \mu \left(\frac{z d^2 \Omega_y}{dy^2} + \frac{S(z) d \Omega_y}{dy} \right) \right]
$$
(14)

Substituting Equation (1) , (2) , (4) and (5) into Equation (10) , we have:

$$
\sigma_y = \frac{E}{1 - \mu^2} \left[\left(-\frac{z d^2}{dy^2} + \frac{S(z) d \cap_x}{dx} \right) - \mu \left(\frac{z d^2}{dx^2} + \frac{S(z) d \cap_y}{dy} \right) \right]
$$
(15)

Substituting Equation (1) , (2) and (6) into Equation (11) , we have:

$$
\tau_{xy} = \frac{E(1-\mu)}{(1-\mu^2)} \left[-\frac{z\partial^2 U}{\partial x \partial y} + S(z) \left(\frac{d \Omega_x}{dy} + \frac{d \Omega_y}{dx} \right) \right]
$$
\nSubstituting Equation (1), (2) and (7) into Equation (12).

Substituting Equation (1) , (2) and (7) into Equation (12) , we have:

$$
\tau_{xz} = \frac{E(1-\mu)}{(1-\mu^2)} \left[\frac{z\partial^2 U}{\partial x \partial z} + S(z) \left(\frac{d \Omega_x}{dz} + \frac{d \Omega_z}{dx} \right) \right]
$$
(17)
Substituting Equation (1), (2) and (8) into Equation (13).

Substituting Equation (1), (2) and (8) into Equation (13), we have:

$$
\tau_{yz} = \frac{E(1-\mu)}{(1-\mu^2)} \left[\frac{z\partial^2 U}{\partial y \partial z} + S(z) \left(\frac{d \Omega_y}{dz} + \frac{d \Omega_z}{dy} \right) \right]
$$
(18)

where:

the symbol ∪ denotes out-of-plane displacement the symbol $F(z)$ denotes Shear deformation profile

the symbol \cap_x denotes shear deformation rotation along x axis the symbol \cap_{ν} denotes shear deformation rotation along y axis the symbol μ denotes poison ratio

the symbol E denotes modulus of elasticity of the plate

the symbol σ_x denotes stress normal to x axis

the symbol σ_{ν} denotes stress normal to y axis

the symbol τ_{xy} denotes shear stress along x-y axis the symbol τ_{xz} denotes shear stress along x-z axis the symbol τ_{vz} denotes shear stress along y-z axis the symbol ε_x denotes normal strain along x axis the symbol ε_{ν} denotes normal strain along y axis the symbol γ_{xy} denotes shear strain along x-y axis the symbol γ_{yz} denotes shear strain along x-z axis

the symbol γ_{vz} denotes shear strain along y-z axis

3.3. Total Potential Energy Equation

The total potential energy expression (\nexists) , was formulated in accordance to the kinematics and constitutive relation in the previous section (see [27]).

$$
\vec{A} = \Delta + \nabla \tag{19}
$$
Where;

$$
\nabla = -\int_0^a \int_0^b w \cup (x, y) \partial x \partial y \tag{20}
$$

 $\frac{1}{20}$ $\frac{1}{20}$ w is the uniformly distributed load.

$$
\Delta = \frac{1}{2} \iiint_{-\frac{t}{2}}^{\frac{t}{2}} (\sigma_x \varepsilon_{x+} \sigma_y \varepsilon_{y+} \tau_{xy} \gamma_{xy} + \tau_{xz} \gamma_{xz} + \tau_{yz} \gamma_{yz}) dx dy dz
$$
 (21)

Thus: $\vec{A} = D \int_0^u \int_0^v \left| \frac{g_1}{2} \right|$ ∂^2 ∪ $\overline{\partial x^2}$ 2 $-g_2$ $\left\lfloor \frac{g_2}{g_1} \right\rfloor$ ∂^2 ∪ $\left(\frac{\partial^2 U}{\partial x^2} \cdot \frac{\partial^2 \Omega_x}{\partial x}\right) + \frac{g_3}{2} \left(\frac{\partial^2 \Omega_x}{\partial x}\right)$ 2 $\bigg\}$ \boldsymbol{b} $\boldsymbol{0}$ \boldsymbol{a} $\boldsymbol{0}$ $+$ | g_1 | ∂^2 ∪ $\overline{\partial x \partial y}$ 2 $-g_2$ $\left\lfloor \frac{g_2}{g_1} \right\rfloor$ ∂^2 ∪ $\frac{\partial^2 U}{\partial x \partial y} \cdot \frac{\partial \Omega_x}{\partial y}$ $-\mathit{g}_{2}\left(\frac{\partial^{2}w}{\partial x\partial y}.\frac{\partial\Omega_{y}}{\partial x}\right)$ $+\left|(1+\mu)\frac{g_3}{2}\left(\frac{\partial\,\Omega_x}{\partial y}\right)\left(\frac{\partial\,\Omega_y}{\partial x}\right)\right|$ $+\frac{(1-\mu)}{4}\bigg|g_3\bigg(\frac{\partial \Omega_x}{\partial y}\bigg)$ 2 $+ g_3 \left(\frac{\partial \Omega_y}{\partial x} \right)$ 2 � $+\left|\frac{g_1}{2}\right|$ ∂^2 ∪ $\overline{\partial y^2}$ 2 $-g_2$ $\left($ ∂^2 ∪ $\frac{\partial^2 U}{\partial y^2} \cdot \frac{\partial \cap_y}{\partial y}$ $+\frac{g_3}{2}\left(\frac{\partial \,\Omega_y}{\partial y}\right)$ 2 � $+\left|\frac{(1-\mu)}{4}g_4(\cap_x)^2\right|$ $+\frac{(1-\mu)}{4}g_4(\cap_y)$ 2 \prod dxd

$$
-\int_0^a \int_0^b \frac{w}{2} \cup (x, y) \partial x \partial y \tag{22}
$$

The non-dimensional values of quantities along the x and yaxis respectively is presented below.

$$
z = ts; x = a \ni \text{ and } y = b \in
$$
 (23)
where;

a, b and t are the length, breath and thickness of the plate while ∋, ∈ and s are the non-dimensional value of length, breath and thickness of the plate.

Let the length to breadth aspect ratio, $\alpha = \frac{b}{b}$ (24)

Span to thickness ratio, $\beta = \frac{a}{t}$ $\frac{1}{t}$ (25)

Deflection (∪), is the product of shape function of the plate and deflection coefficient:

$$
U = C \cdot n \tag{26}
$$

Similarly:

The shear deformation rotation along x-axis becomes:

$$
\cap_x = \left[\frac{dn}{d\ \ 3}\right][C_x]
$$
\nSimilarly, the shear deformation rotation along y-axis

becomes:

$$
\cap_{y} = \left[\frac{an}{d \in}\right] \left[\mathbf{C}_{y}\right] \tag{28}
$$

where,

 n is the shape function of the plate

∁ is the coefficient of deflection

 C_x is the coefficient of shear deformation along x axis

 C_y is the coefficient of shear deformation along y axis

By substituting Equation 23, 24, 25, 26, 27 and 28 into 22, gives:

[∄]⁼ ²⁴ � � ��1∁2 � 2 ∋²� 2 − 22∁∁ � 2 ∋²� 2 1 0 1 0 + 3∁ 2 � 2 ∋²� 2 � + �21 ∁2 [∝]² � 2 ∋ ∈� 2 − 22 ∁∁ ∝2 � 2 ∋ ∈� 2 − 22 ∁∁ ∝2 � 2 ∋ ∈� 2 � + �(1 +)3 ∁∁ [∝]² � 2ℎ ∋ ∈� 2 � + (1 −) ² �3 ∁ 2 ∝2 � 2 ∋ ∈� 2 + 3 ∁ 2 [∝]² � 2 ∋ ∈� 2 � + �1 ∁2 [∝]⁴ � 2 ∈²� 2 − 22 ∁∁ ∝4 � 2ℎ ∈²� 2 + 3 ∁ 2 ∝4 � 2ℎ ∈²� 2 � + � (1 −) ² 24∁ 2 � ∋ � 2 ⁺ (1 −) ² . 24∁ 2 [∝]² � ∈ � 2 �� − � � ∁ ∋ ∈ 1 0 1 0 (29) Given that is the rigidity; let: ⁼³ 12(1 − ²) (30)

3.4. Governing Energy Equation

The elastic plate presented plate in figures 1 and 2 under bending subjected to uniformly distributed load was used to obtain the displacement – strain relationships in terms of curvatures.

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3.4.1. General Governing Equation

The total potential energy shall be minimized with respect to the deflection w, shear deformation along the x axis, \cap_x and shear deformation along y axis∩_y. Minimizing or differentiating total potential energy equation with respect to ∪, \cap_x , and \cap_y is said to be the direct variation.

$$
\frac{\partial \Pi}{\partial U} = \frac{\partial \Pi}{\partial \Omega_x} = \frac{\partial \Pi}{\partial \Omega_y} = 0
$$
\n(31)

By solving the resulting three simultaneous governing equation the actual deflection ∪, shear deformation along the x axis, \cap_x and shear deformation along the y axis, \cap_y was gotten as:

$$
U = \left(a_0 + a_1 \ni + \frac{a_2 \ni^2}{2} + \frac{a_3 \ni^3}{6} + \frac{wa^4}{D} \left(\frac{n_1}{U_3}\right) \cdot \frac{\ni^4}{24}\right)
$$

$$
\times \left(b_0 + b_1
$$

$$
\in + \frac{b_2 \in^2}{2} + \frac{b_3 \in^3}{6} + \frac{wa^4}{D} \left(\frac{n_1}{U_3}\right) \cdot \frac{\in^4}{24}
$$

$$
+ \frac{b_5 \in^5}{120}\right) (32)
$$

The general polynomial shear deformation function (Rotation equation for the y-axis) of a rectangular plate as was obtained and presented:

$$
\begin{aligned} \n\Omega_x &= \left(a_4 + a_5 \ni + \frac{a_6}{2} + \frac{wa^3}{D} \left(\frac{n_4}{g_2 \cap 3} \right) \cdot \frac{\ni^3}{6} \right) \\ \n&\times \left(b_7 + b_8 \right) \\ \n&\in + \frac{b_9 \in^2}{2} + \frac{b_{10} \in^3}{6} + \frac{b_{11} \in^4}{24} \right) \n\end{aligned} \tag{33}
$$

 The general polynomial shear deformation function (Rotation equation for the y-axis) of a rectangular plate as was obtained and presented:

$$
\begin{aligned} \n\Omega_y &= \left(a_7 + a_8 \ni + \frac{a_9}{2} \ni + \frac{a_{10} \ni^3}{6} + \frac{a_{11} \ni^4}{24} \right) \\ \n&\times \left(b_4 + b_5 \right) \\ \n&\in + \frac{b_6 \in^2}{2} + \frac{wa^3}{D} \left(\frac{\infty^3}{g_2 \cap 1} \right) \cdot \frac{\in^3}{6} \right) \tag{34} \n\end{aligned}
$$

 The support conditions of the plate consideration is presented in the Figure 2.

Figure 2: CCFC Rectangular Plate

3.4.2. Direct Governing Equation

 The direct variation method was used to obtain the direct governing differential equation by differentiating the total potential energy with respect to the coefficient of deflection (C), coefficient of shear deformation with respect to x-axis (C_x) and coefficient of shear deformation with respect to y-axis (C_y) . Minimizing the Energy expression (∄) with respect to the coefficient of deflection (∁), coefficient of shear deformation with respect to x-axis (C_x) and coefficient of shear deformation with respect to y-axis (C_{γ}) :

$$
\frac{\partial \vec{\mathbf{\mu}}}{\partial \mathbf{C}} = \frac{\partial \vec{\mathbf{\mu}}}{\partial \mathbf{C}_x} = \frac{\partial \vec{\mathbf{\mu}}}{\partial \mathbf{C}_y} = 0
$$
\n(35)

 The solution of Equation (35) yields the three Equations of equilibrium which is presented as:

$$
t_{11}C - t_{12}C_x - t_{13}C_y = \frac{wa^4}{D} s_q
$$
 (36)

$$
-t_{21}C + t_{22}C_x + t_{23}C_y = 0
$$
\n
$$
-t_{11}C + t_{22}C_x + t_{23}C_y = 0
$$
\n(37)

$$
-t_{31}C + t_{32}C_x + t_{33}C_y = 0
$$
 (38)
Where;

$$
t_{11} = g_1 \left(s_1 + \frac{2}{\alpha^2} s_2 + \frac{1}{\alpha^4} s_3 \right) \tag{39}
$$

$$
t_{12} = -g_2 \left(s_1 + \frac{1}{\alpha^2} s_2 \right) \tag{40}
$$

$$
t_{13} = -g_2 \left(\frac{1}{\alpha^2} s_2 + \frac{1}{\alpha^4} s_3 \right) \tag{41}
$$

$$
t_{21} = -g_2 \left(s_1 + \frac{1}{\alpha^2} s_2 \right) \tag{42}
$$
\n
$$
\left(1 - \mu \right) \tag{42}
$$

$$
t_{22} = \left(g_3 s_1 + \frac{(1-\mu)}{2 \alpha^2} g_3 s_2 + \frac{(1-\mu)}{2} \beta^2 g_4 s_4 \right) \tag{43}
$$

$$
t_{23} = g_3 \frac{(1+\mu)}{2 \alpha^2} s_2 \tag{44}
$$

$$
t_{31} = -g_2 \left(\frac{1}{\alpha^2} s_2 + \frac{1}{\alpha^4} s_3 \right) \tag{45}
$$

$$
t_{32} = g_3 \frac{(1+\mu)}{2 \alpha^2} s_2 \tag{46}
$$

$$
t_{33} = \left(g_3 \frac{(1-\mu)}{2} \left(\frac{1}{\alpha^2} s_2 + \frac{1}{\alpha^4} s_3 \right) + g_4 \frac{(1-\mu)}{2 \alpha^2} \beta^2 s_5 \right) (47)
$$

Given that s_1 , s_2 , s_3 , s_4 , s_5 and s_7 are the stiffness coefficients, let:

$$
s_1 = \int_0^1 \int_0^1 \left(\frac{d^2 n}{d \; \exists^2}\right)^2 d \; \exists \; d \in \tag{48}
$$

$$
s_2 = \int_0^1 \int_0^1 \left(\frac{d^2 n}{d \ni d \in \right)^2} d \ni d \in \tag{49}
$$

$$
s_3 = \int_0^1 \int_0^1 \left(\frac{d^2 n}{d \epsilon^2}\right)^2 d \ni d \epsilon \tag{50}
$$

$$
s_4 = \int_0^1 \int_0^1 \left(\frac{dn}{d\theta}\right)^2 d\theta d\theta \in (51)
$$

$$
s_5 = \int_0^1 \int_0^1 \left(\frac{dn}{d\epsilon}\right)^2 d \ni d\epsilon \tag{52}
$$

$$
s_q = \int_0^{\infty} \int_0^n n \, d \ni d \in \mathfrak{g} \tag{53}
$$

 $\frac{50}{9}$ $\frac{50}{9}$ Equation (36), (37) and (38) simultaneously, gives Equation (54), (55) and (56):

$$
C = \frac{wa^4}{D}(k) \tag{54}
$$

$$
E = \frac{1}{D} (k) \tag{54}
$$

$$
C_x = C M_2 \tag{55}
$$
\n
$$
C_y = C M_3 \tag{56}
$$

$$
Where;
$$

$$
M_2 = \frac{t_{21} \cdot t_{33} - t_{23} \cdot t_{31}}{t_{22} \cdot t_{33} - t_{23} \cdot t_{32}}
$$
(57)

$$
M_3 = \frac{t_{21} \cdot t_{32} - t_{22} \cdot t_{31}}{t_{23} \cdot t_{32} - t_{22} \cdot t_{33}}
$$
(58)

$$
k = \frac{s_6}{t_{11}M_1 - t_{12}M_2 - t_{13}M_3} \tag{59}
$$

3.5. Displacement, Stresses, Bending moment and Shear force Equations

The equations for the moment $(M_x \text{ and } M_y)$, shear force $(Q_x$ and $Q_y)$, in-plane displacement (*u* and *v*), deflection (∪) and the stress of isotropic rectangular thick plate was derived according to [28] by substituting the Equation (54), (55) and (56) where appropriate.

Substituting Equation (54) into (26), gave:

$$
U = \overline{C}n \left(\frac{wa^4}{D}\right)
$$

Let:

$$
\overline{C} = k
$$
 (60)

The bending moment along x-axes:

$$
M_x = \left(-\frac{D_1}{D}wa^2 \overline{C} \left[\frac{d^2 n}{d \overline{\theta}^2} + \mu \frac{d^2 n}{d \overline{\theta}^2} \right] + g_2 \left[\overline{C_x} \frac{d^2 n}{d \overline{\theta}^2} + \mu \overline{C_y} \frac{d^2 n}{d \overline{\theta}^2} \right] \right)
$$
(61)

That is:
\n
$$
M_x = \overline{M_x}wa^2
$$
\n(62)

Given that $D_1 = g_1 D$; and $D_2 = g_2 D$; let: $q_1 = 1$ and $q_2 = 1.2$

$$
\overline{M_x} = -g_1 \overline{C} \left[\frac{d^2 n}{d \overline{\partial}^2} + \mu \frac{d^2 n}{d \overline{\partial}^2} \right] + g_2 \left[\overline{C_x} \frac{d^2 n}{d \overline{\partial}^2} + \mu \overline{C_y} \frac{d^2 n}{d \overline{\partial}^2} \right]
$$
(63)

where;
\n
$$
\overline{C_x} = M_2 \overline{C}
$$
\n(64)

$$
\overline{\mathcal{C}_y} = M_3 \overline{\mathcal{C}}
$$
 (65)

The bending moment along y-axes:

$$
M_{y} = \left(-\frac{D_{1}}{D}wa^{2}C\left[\frac{d^{2}n}{d\epsilon^{2}} + \mu\frac{d^{2}n}{d\epsilon^{2}}\right] + g_{2}\left[\overline{C_{y}}\frac{d^{2}n}{d\epsilon^{2}} + \mu\overline{C_{x}}\frac{d^{2}n}{d\epsilon^{2}}\right]\right)wa^{2}
$$
 (66)

$$
M_y = \overline{M_y}wa^2
$$
 (67)
where;

$$
\overline{M_y} = -g_1 C \left[\frac{d^2 n}{d \epsilon^2} + \mu \frac{d^2 n}{d \epsilon^2} \right] + g_2 \left[\overline{C_y} \frac{d^2 n}{d \epsilon^2} + \mu \overline{C_x} \frac{d^2 n}{d \epsilon^2} \right]
$$
(68)

Shear force
$$
Q_x = \frac{dM_x}{dx}
$$
 and $Q_y = \frac{dM_y}{dy}$ (69)

Thus, the shear force along x-axes:
\n
$$
Q_x = \left(-\overline{C} \left[\frac{\partial^3 n}{\partial \theta^3} + \mu \frac{\partial^3 n}{\partial \theta^3} \right] + \left[\overline{C_x} \frac{\partial^3 n}{\partial \theta^3} + \mu \overline{C_y} \frac{\partial^3 n}{\partial \theta^3} \right] \right) wa
$$
\n(70)

That is:

$$
Q_x = \overline{Q_x}wa
$$
Where; (71)

$$
\overline{Q_x} = -\frac{D_1}{a^3} \overline{C} \left[\frac{\partial^3 n}{\partial \overline{\partial}^3} + \mu \frac{\partial^3 n}{\partial \overline{\partial}^3} \right] + \left[\overline{C_x} \frac{\partial^3 n}{\partial \overline{\partial}^3} + \mu \overline{C_y} \frac{\partial^3 n}{\partial \overline{\partial}^3} \right]
$$
(72)

The shear force along y-axes:

$$
Q_{y} = \left(-\bar{C} \left[\frac{\partial^{3} n}{\partial \theta^{3}} + \mu \frac{\partial^{3} n}{\partial \theta^{3}} \right] + \left[\overline{C_{x}} \frac{\partial^{3} n}{\partial \theta^{3}} + \mu \overline{C_{y}} \frac{\partial^{3} n}{\partial \theta^{3}} \right] \right) wa
$$
\nThat is:

\n
$$
(73)
$$

That is:
\n
$$
Q_y = \overline{Q_y}wa
$$
\n(74)

where;

$$
\overline{Q_y} = -\frac{D_1}{a^3} \overline{C} \left[\frac{\partial^3 n}{\partial \overline{\partial}^3} + \mu \frac{\partial^3 n}{\partial \overline{\partial}^3} \right] + \left[\overline{C_x} \frac{\partial^3 n}{\partial \overline{\partial}^3} + \mu \overline{C_y} \frac{\partial^3 n}{\partial \overline{\partial}^3} \right] \tag{75}
$$

The in-plane displacement along x-axes:

$$
u = \left[-\overline{C}s + \overline{C_x}S(s) \right] \frac{dn}{d \theta} \left(\frac{wa^4}{\beta D} \right) \tag{76}
$$

The in-plane displacement along y-axes:

$$
v = \frac{1}{\alpha} \left[-\overline{\zeta}s + \overline{\zeta_y}S(s) \right] \frac{dn}{d\epsilon} \left(\frac{twa^3}{D} \right)
$$
 (77)
The normal stress along x-aves:

The normal stress along x-axes:

$$
\sigma_x = 12 \left[[-\overline{\zeta}s + \overline{\zeta_x}S(s)] \frac{d^2n}{d\overline{\theta}^2} + \frac{\mu}{\alpha^2} [-\overline{\zeta}s + \overline{\zeta_y}S(s)] \frac{d^2n}{d\overline{\theta}^2} \right] (w\beta^2)
$$
 (78)
The normal stress along v-aves:

he normal stress along y-axes:

$$
\sigma_y = w\beta^2 \left[12 \left[\mu \left[-\overline{Cs} + C_x S(s) \right] \frac{d^2 n}{d \theta^2} \right] + \frac{\mu}{\alpha^2} \left[-\overline{Cs} + \overline{C_y} S(s) \right] \frac{d^2 n}{d \theta^2} \right]
$$
\nThe shear stress along x y axes:

\n
$$
y = \frac{1}{\alpha^2} \left[-\overline{C_x} + \overline{C_y} S(s) \right] \frac{d^2 n}{d \theta^2} \right]
$$

 The shear stress along x-y axes: $\langle \cdot, \cdot \rangle$ r

$$
\tau_{xy} = 6 \frac{(1 - \mu)}{\alpha} \Big| -2 \overline{C} s + \overline{C}_x S(s) + \overline{C}_y S(s) \cdot \frac{1}{\alpha} \Big| \frac{d^2 n}{\partial \theta \partial \theta} (\omega \beta^2)
$$
\n
$$
\text{The shear stress along x-z axes:} \tag{80}
$$

$$
\tau_{xz} = 6(1-\mu)\overline{\mathcal{C}_x} \frac{dS(z)}{dz} \frac{dn}{d\overline{z}}(w\beta^2)
$$
(81)

The shear stress along y-z axes:

$$
\tau_{yz} = \frac{6(1-\mu)}{\alpha} \overline{C}_y \frac{dS(z)}{dz} \frac{dn}{d \epsilon} (w\beta^2)
$$
\n
$$
q_d = \text{Self weight of the plate}
$$
\n(82)

4. Results and Discussions

www.astesj.com 1433 The numerical analysis of a CCFC rectangular plate of various span-thickness ratios is presented in Table 2 to 6. A third-

order polynomial displacement function for the analysis CCFC as was derived in Equation (32) in line with the work of [29] is presented in Equation 83 to 94:

$$
U_{(\exists,\in)} = \left(a_0 + a_1 \ni \frac{a_2 \ni^2}{2} + \frac{a_3 \ni^3}{6} + \frac{\ni^4}{24} Sa_4\right) \cdot \left(b_0 + b_1\right)
$$

$$
\in +\frac{b_2 \in^2}{2} + \frac{b_3 \in^3}{6} + \frac{\in^4}{24} b_4 + \frac{b_5 \in^5}{120}\right) \quad (83)
$$

where;

$$
S = \frac{wa^4}{D} \left(\frac{n_1}{u_3}\right)
$$

At \exists = \in = 0; U = 0
At \exists = \in = 0: Slone (i) $\frac{dU}{d} = \frac{dU}{d} = 0$ (84)

At
$$
\exists
$$
 = \in = 0; *Slope* (*ie*. $\frac{d0}{d\exists} = \frac{d0}{d\in} = 0$) (85)

At
$$
\exists
$$
 = \in = 1; \cup_x = 0 and M_y (*ie*. $\frac{d^2\cup}{d\in^2}$ = 0) (86)

At
$$
\ni = 1
$$
, $slope\left(ie.\frac{d\upsilon}{d\upsilon} = 0\right)$; at $\epsilon = 1$, $slope\left(ie.\frac{d\upsilon}{d\epsilon} = \frac{2}{3b_5}\right)$ (87)

At
$$
\epsilon = 1
$$
; shear force $Q_y \left(i e. \frac{d^3 w}{d \epsilon^3} = 0 \right)$ (88)

Substituting Equations (84 to 88) into Equation (83) and solving gives the following constants:

$$
a_0 = 0; a_1 = 0; a_2 = \frac{S_{a4}}{12}; a_3 = \frac{-S_{a4}}{2} \text{ and } b_0 = 0; b_1
$$

= 0; b_2 = 2.8b_5; b_3 = -5.2b_5; S_{b4}
= 3.8b_5 \t(89)

Substituting the constants of Equation (89) into Equation (83) gives;

$$
U = \frac{S_{a4}}{24} (3^2 - 2 \ 3^3 + 3^4)
$$

$$
\times \frac{b_5}{360} (2.8 \ e^2 - 5.2 \ e^3 + 3.8 \ e^4 - e^5) \quad (90)
$$

That is:

 $U = \frac{S_{a4} \times b_5}{8640}$ (∋² – 2 ∋³ + ∋⁴)

$$
\times
$$
 (2.8 \in ² - 5.2 \in ³ + 3.8 \in ⁴ - \in ⁵) (91)
Recall from Equation 26, that;

$$
U = n.C
$$

Let the amplitude:

$$
C = \frac{1}{8640} (S_{a4} \times b_5)
$$

$$
\frac{1}{540}(S_{a4} \times b_5)
$$

$$
C = \frac{1}{8640} (S_{a4} \times b_5)
$$
\nand;
\n
$$
n = (3^2 - 2 \cdot 3^3 + 3^4) \times (2.8 \cdot 5^2 - 5.2 \cdot 5^3 + 3.8 \cdot 5^4 - 5^5) (93)
$$

Therefore:
\n
$$
U = \frac{S_{a4}.b_5}{8640} (1.5 \text{ }^32 - 2.5 \text{ }^34 + 3^4)
$$
\n
$$
\times (2.8 \text{ }^22 - 5.2 \text{ }^34 + 3.8 \text{ }^42 - 5^5) \tag{94}
$$

The values stiffness coefficient obtained from the above expression is presented in Table 1. Table 2 contains the result of bending moment, shear force, and their resultants of a square CCFC rectangular plate at different span to thickness aspect ratio. These numerical values of the design factors of these quantities $(\overline{U}, \overline{M}$ and \overline{Q}) for a square thick plate is obtained using Equation (60) to (75) as obtained from this study. Table 3 contains the result of the values of displacement (u, v and ∪) and the stresses characteristics $(\sigma_x, \sigma_y, \tau_{xy}, \tau_{xz} \text{ and } \tau_{yz})$ of a square CCFC

rectangular plate at different span to thickness aspect ratio using Equation (60), (70) and (77). These numerical values were obtained from the Equation (60), (78) to (82). Table 4 contains the result of the value of displacement (u, v and ∪) and the stresses characteristics $(\sigma_x, \sigma_y, \tau_{xy}, \tau_{xz} \text{ and } \tau_{yz})$ for length to breadth ratio of 1.5 for CCFC rectangular plate at different span to thickness aspect ratio. These numerical values were obtained from the Equation (60), (78) to (82). Table 5 contains the result of the values of displacement (u, v and ∪) and the stresses characteristics $(\sigma_x, \sigma_y, \tau_{xy}, \tau_{xz}$ and τ_{yz}) for length to breadth ratio of 2.0 for CCFC rectangular plate at different span to thickness aspect ratio. These numerical values were obtained from the Equation (60), (78) to (82). A comparative analysis of results obtained from this work is presented in Table 6 and Figure 3. This is determined the correctness of the results from the present studies, comparison was made between values from the present study and those of past scholars. Table 2 presents the result of the comparison done with the work of [30] for CCFC rectangular thick plate of 2.0 lengths to breadth ratio of various aspect ratios and the corresponding percentage difference between the values of centroidal deflection. Figure 3 contains a graph of the deflection versus span to the thickness aspect ratio of CCFC thick plate of length to breadth ratio of 2.0. This shows the comparison between the present work and that of the author in [30].

4.1. Discussion

The numerical results of stiffness coefficient of the plate is obtained from Equation (31) to (36).

Table 1: **S**tiffness coefficient, s values for CCFC boundary conditions

Theor	Plate	s_{1}	S_{2}	s_{2}	S_A	$S_{\rm E}$	$S_{\mathcal{L}}$
Presen (NPST	CCFC	0.05	0.00	0.00	0.00	0.000	0.00

Table 2: Bending Moments, Shear Force and Stress resultants of CCFC plate for $b/a = 1.0$

It is seen from Table 2 that as the deflection (∪), bending moment $(M_x \text{ and } M_y)$ and $(Q_x \text{ and } Q_y)$ decreases, the span to thickness ratio increases. These continue until failure occurs in the plate structure. This means that, the load that causes the plate to deflect also causes the plate material to bend simultaneously. It is observed that the value of deflection varies less as the span to thickness increase, this equal to the value of the CPT at span to thickness ratio of 100 and above. Table 2 reveals that the values of deflection at both the x and y axis decrease as the spanthickness ratio increases with a constant value of 0.0092 at the span-thickness ratio of 30.

Table 3: Displacement and Stresses of CCFC plate for length to breadth ratio of 1.0

	U	\boldsymbol{u}	$\sigma_{\!x}$	τ_{xy}	τ_{yz}
β	Ū	ū	$\overline{\sigma_{x}}$	$\overline{\tau_{xy}}$	$\overline{\tau_{yz}}$
4	0.0165	-0.0038	0.2562	-0.0635	0.0229
10	0.0103	-0.0030	0.2000	-0.0479	0.0031
15	0.0096	-0.0029	0.1943	-0.0462	0.0014
20	0.0094	-0.0029	0.1924	-0.0457	0.0008
25	0.0093	-0.0028	0.1915	-0.0454	0.0005
30	0.0093	-0.0028	0.1910	-0.0453	0.0004
35	0.0092	-0.0028	0.1907	-0.0452	0.0003
40	0.0092	-0.0028	0.1905	-0.0451	0.0002
45	0.0092	-0.0028	0.1904	-0.0451	0.0002
50	0.0092	-0.0028	0.1903	-0.0451	0.0001
55	0.0092	-0.0028	0.1902	-0.0450	0.0001

60	0.0092	-0.0028	0.1902	-0.0450	0.0001
65	0.0092	-0.0028	0.1901	-0.0450	0.0001
70	0.0092	-0.0028	0.1901	-0.0450	0.0001
75	0.0092	-0.0028	0.1900	-0.0450	0.0001
80	0.0092	-0.0028	0.1900	-0.0450	0.0001
85	0.0092	-0.0028	0.1900	-0.0450	0.0000
90	0.0092	-0.0028	0.1900	-0.0450	0.0000
95	0.0092	-0.0028	0.1900	-0.0450	0.0000
100	0.0092	-0.0028	0.1900	-0.0450	0.0000
1000	0.0092	-0.0028	0.1900	-0.0450	0.0000
CPT	0.0092	-0.0028	0.1900	-0.0450	0.0000

From the results of SCFS plate presented in Figure 12 to 20, it shows that the values of critical lateral imposed load q_{iw} decrease as the length-width ratio increases, this continues until failure occurs. This means that an increase in plate length increases the chance of failure in a plate structure. Meanwhile the values of critical lateral imposed load q_{ip} increase as the lengthwidth ratio increases, this continues until safety is ensured.

From a critical look at Table 3, it is shown that the value of displacement (u, v and ∪) characteristics decreases with increases in the value of the span-thickness ratio. It is also observed in the Tables that the displacement (u, v, w) and v) and stresses characteristics increases as the value of the length to breadth ratio increases. This means that, the in-plane displacement is functions of x, y and z as it varies with the plate thickness while the deflection is only a function of x and y and did not vary linearly with the thickness of the plate thickness. Also, it was deduced that the normal stress (σ_x and σ_y) and shear stress characteristics $(\tau_{xy}, \tau_{xz} \text{ and } \tau_{yz})$ also decreases as the span-thickness ratio increases. It is also observed in the Table 3 that the stresses characteristics $(\sigma_x, \sigma_y, \tau_{xy}, \tau_{xz} \text{ and } \tau_{yz})$ increases as the value of the length to breadth ratio increases.

It can be seen that, at span to thickness ratio between 4 and 15, the value of transverse shear stress along y and z axes (τ_{vz}) varies between 0.022904 and 0.001370. These values of transverse shear stress (τ_{yz}) decrease between 0.000765 and 0.000100 at the span to thickness between 20 and 55 respectively. Hence, the value of transverse shear stress (τ_{vz}) is about 0.000084 and 0.00003, at the span to thickness between 60 and 100. The value becomes almost constant or equal to the value from CPT. This value becomes negligible as the value span to thickness increases. The values of transverse stress for spanthickness ratios of 100 and above are equal to the value from CPT.

In summary, there are three categories of rectangular plates. The plates whose vertical shear stress do not vary well from zero will be classified as thin plates because its value is almost equal to the value of the CPT. In between the thin and thick plate is the classified as moderate thick plate. Since the plate whose transverse shear stress varies very much from zero is categorized as thick plates. Therefore, the span-to-depth ratio for these categories of rectangular plates are: Thick plate: $a/t \leq 15$; moderately thick plate: $20 \le a/t \le 55$; thin plate: $a/t \ge 60$. This confirmation can be used to show the boundary between thin and thick plate. Thus, it can be deduced from this research work that thick plate is the one whose span-depth ratio value is 4 up to 15.

Table 4: Displacement and Stresses of CCFC plate for length to breadth ratio of 1.5

	U	$\boldsymbol{\mathcal{u}}$	σ_x	τ_{xy}	τ_{yz}
β	Ū	ū	$\overline{\sigma_{\rm x}}$	$\overline{\tau_{xy}}$	$\overline{\tau_{yz}}$
4	0.0185	-0.0043	0.2768	-0.0491	0.0126
10	0.0115	-0.0033	0.2157	-0.0361	0.0017
15	0.0108	-0.0032	0.2096	-0.0348	0.0007
20	0.0106	-0.0032	0.2075	-0.0343	0.0004
25	0.0105	-0.0032	0.2065	-0.0341	0.0003
30	0.0104	-0.0032	0.2059	-0.0340	0.0002
35	0.0104	-0.0032	0.2056	-0.0339	0.0001
40	0.0104	-0.0032	0.2054	-0.0338	0.0001
45	0.0104	-0.0032	0.2053	-0.0338	0.0001
50	0.0103	-0.0032	0.2052	-0.0338	0.0001
55	0.0103	-0.0032	0.2051	-0.0338	0.0001
60	0.0103	-0.0032	0.2050	-0.0338	0.0000
65	0.0103	-0.0032	0.2050	-0.0338	0.0000
70	0.0103	-0.0032	0.2050	-0.0338	0.0000
75	0.0103	-0.0032	0.2050	-0.0337	0.0000
80	0.0103	-0.0032	0.2049	-0.0337	0.0000
85	0.0103	-0.0032	0.2049	-0.0337	0.0000
90	0.0103	-0.0032	0.2049	-0.0337	0.0000
95	0.0103	-0.0032	0.2049	-0.0337	0.0000
100	0.0103	-0.0032	0.2048	-0.0337	0.0000
1000	0.0103	-0.0032	0.2048	-0.0337	0.0000
CPT	0.0103	-0.0032	0.2048	-0.0337	0.0000

From a critical look at Table 4, it is shown that the nondimensional displacement (u, v and ∪) characteristics decrease with increases in the value of the span-thickness ratio. It is also observed in the Tables that the displacement (u, v and ∪) and stresses characteristics increase as the value of the length to breadth ratio increases. This means that, the in-plane

displacement are functions of x, y and z as it vary with the plate thickness while the deflection is only a function of x and y and did not varies linearly with the thickness of the plate thickness. Also, it was deduced that the normal stress $(\sigma_x$ and σ_y) and shear stress characteristics $(\tau_{xy}, \tau_{xz} \text{ and } \tau_{yz})$ also decrease as the spanthickness ratio increases. It is also observed in the Table 4 that the stresses characteristics $(\sigma_x, \sigma_y, \tau_{xy}, \tau_{xz} \text{ and } \tau_{yz})$ increase as the value of the length to breadth ratio increases.

It can be seen that, at span to thickness ratio between 4 and 10, the value of transverse shear stress along y and z axes (τ_{yz}) varies between 0.012564 and 0.001667. These values of transverse shear stress (τ_{vz}) decrease between 0.000724 and 0.000100 at the span to thickness between 15 and 40. Hence, the value of transverse shear stress (τ_{yz}) is about 0.000079 and 0.000016, at the span to thickness between 45 and 100. The value becomes almost constant or equal to the value from CPT. This value becomes negligible as the value span to thickness increases. The values of non-dimensional value of transverse stress for spanthickness ratios of 100 and above are equal to the value from CPT.

In summary, there are three categories of rectangular plates. The plates whose vertical shear stress does not very well from zero will be classified as thin plates because its value are almost equal to the value of the CPT. In between the thin and thick plate is the classified as moderate thick plate. Since the plate whose transverse shear stress varies very much from zero is categorized as thick plates. Therefore, the span-to-depth ratio for these categories of rectangular plates are: Thick plate: $a/t \leq 10$; moderately thick plate: $15 \le a/t \le 40$; thin plate: $a/t \ge 45$. This confirmation can be used to show the boundary between thin and thick plate. Thus, it can be deduced from this research work that a thick plate is the one whose span-depth ratio value is 4 up to 10.

Table 5: Displacement and Stresses of CCFC plate for length to breadth ratio of 2.0

	U	\boldsymbol{u}	σ_x	τ_{xy}	τ_{yz}
β	ū	ū	$\overline{\sigma_{\rm x}}$	$\overline{\tau_{xy}}$	$\overline{\tau_{yz}}$
4	0.0191	-0.0044	0.2824	-0.0386	0.0083
10	0.0119	-0.0034	0.2198	-0.0281	0.0011
15	0.0112	-0.0033	0.2136	-0.0270	0.0005
20	0.0110	-0.0033	0.2114	-0.0266	0.0003
25	0.0109	-0.0033	0.2104	-0.0265	0.0001
30	0.0108	-0.0033	0.2098	-0.0264	0.0001
35	0.0108	-0.0033	0.2095	-0.0263	0.0001
40	0.0107	-0.0033	0.2093	-0.0263	0.0001
45	0.0107	-0.0033	0.2092	-0.0262	0.0001
50	0.0107	-0.0033	0.2091	-0.0262	0.0000

From a critical look at Table 5, it is shown that the nondimensional displacement (u, v and ∪) characteristics decrease with increases in the value of the span-thickness ratio. It is also observed in the Tables that the displacement (u, v and ∪) and stresses characteristics increase as the value of the length to breadth ratio increases. This means that, the in-plane displacement are functions of x, y and z as it vary with the plate thickness while the deflection is only a function of x and y and did not varies linearly with the thickness of the plate thickness. Also, it was deduced that the normal stress (σ_x and σ_y) and shear stress characteristics $(\tau_{xy}, \tau_{xz} \text{ and } \tau_{yz})$ also decrease as the spanthickness ratio increases. It is also observed in the Table 4 that the stresses characteristics $(\sigma_x, \sigma_y, \tau_{xy}, \tau_{xz} \text{ and } \tau_{yz})$ increase as the value of the length to breadth ratio increases.

It can be seen that, at span to thickness ratio between 4 and 10, the value of transverse shear stress along y and z axes (τ_{vz}) varies between 0.008325 and 0.001101. These values of transverse shear stress (τ_{yz}) decrease between 0.000478 and 0.000118 at the span to thickness between 15 and 30 respectively. Hence, the value of transverse shear stress (τ_{vz}) is about 0.000087 and 0.000011, at the span to thickness between 35 and 100. The value becomes almost constant or equal to the value from CPT. This value becomes negligible as the value span to thickness increases. The values of non-dimensional value of transverse stress for spanthickness ratios of 100 and above are equal to the value from CPT.

In summary, there are three categories of rectangular plates. The plates whose vertical shear stress do not vary well from zero will be classified as thin plates because its value are almost equal to the value of the CPT. In between the thin and thick plate is the classified as moderate thick plate. Since the plate whose transverse shear stress varies very much from zero is categorized as thick plates. Therefore, the span-to-depth ratio for these categories of rectangular plates are: Thick plate: $a/t \leq 10$; moderately thick plate: $10 \le a/t \le 30$; thin plate: $a/t \ge 35$. This confirmation can be used to show the boundary between thin and thick plate. Thus, it can be deduced from this research work that thick plate is the one whose span-depth ratio value is 4 up to 10.

In the classical plate theory (CPT), it is assumed that the plane cross sections initially normal to the plate's mid-surface before deformation remain plane and normal to the mid-surface after deformation , and the length of such elements is not altered [31]. This is the result of neglecting the transverse shear strains, therefore value of vertical shear stress from CPT analysis is zero. Any plate whose span-to-thickness ratio is such that the value of transverse shear stress from thick plate analysis is approximately zero can be idealized as thin plate. Analyzing such plate with classical plate theory will not introduce significant errors.

Thus, these plates can be classified as moderately thick plate. Hence, analyzing them with classical plate theory will introduce significant errors. When the span-to-thickness ratio is less than 15, the value of vertical shear stress is significant even when corrected to 4 decimal places. This range of span-to-thickness ratio produces plate classified as thick plate. It is also, showed that at span-thickness ratio above 100, the values obtained from the models used herein coincide exactly with values from CPT. This is quite expected since we assumed in CPT analyses that at spanthickness ratios from 100 and above, a plate can be taking as being thin.

Table 6: Comparative analysis of results from the present work and literature for 2.0 aspect ratio (α)

ß	Present	$[30]$	Percentage
	work		difference
		u(m)	(%)
	u(m)		
4	-0.0044	-0.0037	20.772
10	-0.0034	-0.0033	3.3903
50	-0.0033	-0.0033	0.2749
100	-0.0033	-0.0033	0.0306
1000	-0.0033	-0.0033	0.0306
CPT	-0.0033	-0.0033	
Average			
total $%$		4.9	
difference			

From Table 6, it is shown that the present study predicts slightly higher values of in-plane displacement along the direction of x coordinate with highest percentage difference of 20.8% which occur at length to breadth aspect ratio of 4. This proves the safety of the present model. The present study converges with the work as value of span to thickness ratio increases and becomes the same or equal to CPT value at span to thickness ratio of 100. The rate of convergence can be seen in Figure 3. The total average percentage difference between the present study and the author in [30] is 4.9%. This means that at about 95 % confidence level, the values from the present study are the same with those of [30]. This higher confidence level proved that the polynomial shear deformation theory can be used with confidence for stress and bending analysis of thick rectangular plate clamped on three edges, and free at the other remaining edge.

5. Conclusion

Stress analysis of thick rectangular plate using higher order polynomial shear deformation theory has been investigated. From the study this conclusion has been drawn:

- i. It is concluded that the values of transverse shear stress obtained by this theory achieve accepted vertical shear stress to the thickness of plate variation and satisfied the transverse flexibility of condition at the top and bottom faces of the plate while predicting the flexural characteristics for an isotropic rectangular CCFC plate.
- ii. The result moment and stresses obtained by present work using polynomial shear deformation theory agreed well with those of refined plate theory, but varied more with value of CPT. This validates the efficacy of the present theory in the thick plate analysis.
- iii. The governing differential equations and associated boundary conditions obtained are variationally consistent and can be used with confidence in the analysis of rectangular thick plate.

6. Recommendation of Future Scope of Studies

This lies in the establishment shear deformation/displacement function for free vibration and buckling analysis of rectangular or circular thick plate. Then determining the displacement and stress parameters for the plate using virtual work principles.

Figure 4: Graph of span to thickness ratio versus stresses of a square plate

Figure 5: Graph of span to thickness ratio versus stresses for length to breadth ratio of 1.0

Figure 6: Graph of span to thickness ratio versus stresses for length to breadth ratio of 1.5

Figure 7: Graph of span to thickness ratio versus stresses for length to breadth ratio of 2.0

Conflict of Interest Statement

There is no conflict of interest associated with this work.

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